Simple Extension of a Hypo-elastic Body of Grade Zero

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1. **Introduction.** The constitutive equations of an isotropic hypo-elastic body, as defined by Truesdell (1955a), are

$$\tilde{s}^{ik}g_{ki} = h_0 d_k^k \delta_i^i + h_1 d_i^i + h_2 d_k^k s_i^i + h_3 M \delta_i^i
+ \frac{1}{2} h_4 (d_k^i s_i^k + s_k^i d_i^k) + h_5 d_k^k s_r^i s_i^r
+ h_6 M s_i^i + h_7 N \delta_i^i + \frac{1}{2} h_8 (d_r^i s_m^r s_i^m + s_r^i s_m^r d_i^m)
+ h_9 M s_r^i s_i^r + h_{10} N s_i^i + h_{11} N s_k^i s_i^k ,$$
(1.1)

$$d_{ij} = \frac{1}{2}(v_{i,j} + v_{j,i}), \qquad M = s_i^i d_i^i, \qquad N = s_i^i s_k^i d_i^k,$$
 (1.2)

$$\tilde{s}^{ik} = \frac{\partial s^{ik}}{\partial t} + v^m s^{ik}_{,m} - s^{im} v^k_{,m} - s^{mk} v^i_{,m} + s^{ik} d^m_{,m}, \qquad (1.3)$$

where g_{ik} is the covariant metric tensor of a general fixed system of coordinates, v_i is the covariant velocity vector, $2\mu s_i^i$ is the stress tensor, μ being the rigidity modulus of classical elasticity, and where a comma denotes covariant differentiation. Also $h_1, h_2, \dots h_{11}$ are dimensionless analytic functions of the principal invariants of s_i^i . To complete the system of differential equations we add the equations of motion and continuity in the form

$$2\mu s_{,i}^{ij} = \rho \left(\frac{\partial v^{i}}{\partial t} + v_{,i}^{i} v^{j} \right),$$

$$\frac{\partial \rho}{\partial t} + (\rho v^{i})_{,i} = 0,$$
(1.4)

where ρ is the density and body forces are zero.