On the Completeness of the Boussinesq-Papkovich Stress Functions*

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1. Introduction.

In the absence of body forces, the displacement equations of equilibrium in the linear theory of homogeneous and isotropic elastic media, appear as

$$\nabla^2 \mathbf{u} + \frac{1}{1 - 2\nu} \nabla \nabla \cdot \mathbf{u} = 0, \tag{1.1}$$

where \mathbf{u} is the displacement field, while ν denotes Poisson's ratio.

Papkovich [1],² [2] in 1932 established the general solution of (1.1) in the form

$$\mathbf{u} = \nabla(\varphi + \mathbf{r} \cdot \boldsymbol{\psi}) - 4(1 - \nu)\boldsymbol{\psi}, \tag{1.2}$$

where

$$\nabla^2 \varphi = 0, \qquad \nabla^2 \psi = 0. \tag{1.3}$$

Here \mathbf{r} is the position vector of a field point, referred to the origin O. Substitution into (1.1) at once confirms that the vector field (1.2) satisfies the equilibrium equations provided that the generating stress functions (displacement potentials) φ , ψ are sufficiently smooth and meet (1.3). Conversely, every sufficiently regular solution of (1.1) admits the representation (1.2), (1.3), as shown by Parkovich (loc. cit.). A more direct proof of the completeness of this representation is due to Mindlin [3].

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¹Throughout this paper bold face letters designate vectors; the symbols "•" and "★" indicate scalar and vector multiplication of two vectors, respectively.

□ is the usual del-operator.

²Numbers in brackets refer to the bibliography at the end of the paper.

⁸See also Sokolnikoff [4], p. 329 *et seq.*, where the same argument is presented. Mindlin [5] recently extended his proof to the case where the body forces are included; in this instance the Laplace equations (1.3) need to be replaced by appropriate Poisson equations.